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# **The Application of Phase Relationships to Complex Structures III. The Optimum Use of Phase Relationships**

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An improvement is described in the automatic procedure for solving crystal structures incorporated in the computer program *LSAM.* The development of signs from an initial set containing symbols is carried only as far as is necessary to establish strong relationships between the symbols. The information so gained is used in a fresh beginning of the symbolic-addition process. Some failure of relationships between symbols is allowed to give a multisolution method. A phase-permutation computer program for non-centrosymmetric structures, *MULTAN,* incorporates a weighted tangent formula. This is of the form

$$
\tan \varphi_{\mathbf{h}} = \frac{\sum_{\mathbf{h}'} w_{\mathbf{h}'} w_{\mathbf{h} - \mathbf{h}'} |E_{\mathbf{h}'} E_{\mathbf{h} - \mathbf{h}'}| \sin (\varphi_{\mathbf{h}'} + \varphi_{\mathbf{h} - \mathbf{h}'})}{\sum_{\mathbf{h}'} w_{\mathbf{h}'} w_{\mathbf{h} - \mathbf{h}'} |E_{\mathbf{h}'} E_{\mathbf{h} - \mathbf{h}'}| \cos (\varphi_{\mathbf{h}'} + \varphi_{\mathbf{h} - \mathbf{h}'})} = \frac{T_{\mathbf{h}}}{B_{\mathbf{h}}}
$$

and

 $w_h = \tanh{\{\sigma_3 \sigma_2^{-3/2} |E_h| (T_h^2 + B_h^2)^{1/2}\}}$ .

All phases are accepted as soon as they are found with the associated weight. This gives a fourfold increase in speed in development of the complete phase set. An absolute figure of merit is described to indicate probably correct phase sets for multisolution methods.

Direct methods have now advanced to the point where it can be claimed, with some justification, that they are the major crystallographic tool for the solution of equal-atom structures. There are regularly reported the routine solutions of structures which, not many years ago, would have been regarded as veritable tours de force. However the very efficiency of these methods can lead to a lack of sophistication in their use since even the crudest techniques of application produce results comparable with those of quite intricate alternative methods.

The general philosophy of our approach is to stretch the usefulness of direct methods as far as possible. They are not to be regarded just as an easy alternative to solving structures in some other way but rather as having the potentiality of development to the point where they can solve structures which could not otherwise reasonably be solved. It turns out that there are good and bad ways of using any given data set and that with the right techniques structures can be solved which do not yield to the most obvious methods of applying sign and phase relationships. We shall here describe the methods which we have found to be the most effective and to have given solutions when other methods have failed.

# **Centrosymmetric structures**

The multisolution symbolic-addition process for centrosymmetric structures described by Germain & Woolfson (1968) has been fully automated in the *LSAM* (Logical Symbolic Addition Method) computer program. The program goes through the stages:

(I) There are allocated origin-fixing signs to some reflexions and sign symbols to up to six others.

(II) From the  $\sum_{2}$  relationship there is found the reflexion whose sign (perhaps in symbolic form) is given with the highest probability. This is added to the list of known signs and the process repeated until the highest probability is lower than some preset limit,  $P_L$ .

(III) All the triple-product sign relationships involving reflexions in the known list are examined to find relationships between sign symbols.

(IV) All sets of signs for symbols are found which satisfy a specified number of the strongest relationships between sign symbols. Each of these sets constitutes a possible solution.

(V) For each possible set of signs various figures of merit are computed which normally enable the correct solution to be recognized.

Relationships between sign symbols, found at stage III, are related to the situation in stage II when a reflexion is found in terms of two different combinations of symbols. Thus a reflexion may have an indicated sign

*AB* with probability 0.90  
and *CD* with probability 0.85. 
$$
(1)
$$

If 0.90 is greater than  $P<sub>L</sub>$  and is the highest probability for any new sign indication at this stage then the

# Table 1. Computer output for the sign determining process

### 2,3,4,5 TETRAPHENYLCYCLCPENTACIENE

10 JULY 1969



2,3,4,5 TETRAPHENYLCYCLCPENTACIENE

#### 10 JULY 1969

RELATIONSPIPS BETWEEN SYMBOLS



symbol *AB* will be accepted for this reflexion. Subsequently at stage III there will be a relationship between symbols

$$
A B C D \simeq +1.
$$

In practice, for reasonably complex structures, to get most reflexions determined in terms of symbols, *PL*  must be made as low as about 0-80. If several sign indications are accepted at this level of probability there is a non-negligible likelihood that some of them may be wrong.

There is another aspect of this situation which can be seen from the example given in (1). If *AB* and *CD*  are actually similar sign indications, *i.e.* both plus or both minus, with probabilities  $P_1$  and  $P_2$  then the actual probability that the new sign indication is *AB* (or *CD)* is

$$
P = \frac{P_1 P_2}{1 - P_1 - P_2 + 2P_1 P_2}.
$$
 (2)

For the numerical values quoted this gives  $P = 0.98$ , a very high probability. However if *AB* and *CD* are actually opposite indications then the resultant probability for the sign being *AB* is only 0.61. In this event the chance of error is high and if prior knowledge of the signs represented by *A, B, C* and D was available the next reflexion to have its sign accepted would certainly not be this one.

The *LSAM* program in its original form has only once failed to solve a structure in such a way that the best set of signs gave an easily recognizable and substantially correct Fourier map. This was for the structure 2,3,4,5-tetraphenylcyclopentadiene, space group *C2/c* with eight molecules in the unit cell. The *LSAM* program produced seven sets of signs; the map from the set with the highest figures of merit gave a recognizable molecule in a location which was obviously incorrect from packing considerations. The structure was eventually solved by a process which involved relaxing the space group to  $C2$  (Evrard, 1971). When the *LSAM* determined signs were compared with those from the refined structure it was found that about one-third of them had been incorrect.

*A 'post mortem'* on this structure suggested that at some fairly early stage in sign development a wrong indication had been accepted. From this investigation there emerged a revision of the sign-development technique which is incorporated in the latest version of the program, *LSAM 2.* 

In applying *LSAM* 2 to 2,3,4,5-tetraphenylcyclopentadiene, 393 reflexions with  $E > 1.5$  were utilized. The program selected two reflexions to fix the origin (space group *C2/c)* and three others to which it allocated letter symbols  $A$ ,  $B$  and  $C$ . New signs were generated as in stage II but with a high probability limit,  $P_L = 0.95$ . This caused the sign determining process to terminate with 86 signs developed. In Table 1 the computer output is shown at this stage. It is to be interpreted as follows: reflexion 36 has sign *BC* with

probability 0.9999 while reflexion 131 has sign  $-B$ with probability 0.9569. Where the probability of an indication is less than 0.95 that reflexion has not been used in  $\Sigma_2$  relationships to deduce extra signs.

The program next searched for relationships between symbols and found three

$$
AC = -1
$$
, probability 1.0000  

$$
A = +1
$$
, probability 0.9999  
and 
$$
C = -1
$$
, probability 0.9998.

These are shown at the foot of Table 1.

We see from this that to get strong relationships between symbols does not require the development of signs for most of the reflexions and that one can insist on much higher probabilities up to this stage. *LSAM*  next considered the possibilities

$$
A = +1, B = +1, C = -1
$$
  
and 
$$
A = +1, B = -1, C = -1
$$

and started again from the beginning developing signs for each combination. The output from this stage is shown in Table 2. The second set had the higher figures of merit (Germain & Woolfson, 1968) and of the 372 signs determined in this set only 3 were incorrect. In fact *LSAM* was prepared to allow other combinations of sign for A, B and  $C - \text{up to six in}$  $all - but the operator intervened so only the two$ most probable ones were considered.

It is possible to deduce why the change of procedure gave a successful outcome in this case. If it is assumed that the relationships from which the initial 86 signs were determined were all independent then the probability that all these signs were determined without error was the product of the individual probabilities for each sign being correct. This gives  $P_+$  equal to about 0.70 and it is therefore more likely that all the signs are correctly found than that there should be any error. The true value of  $P_+$  is higher since it is known that probabilities given by the usual hyperbolic tangent formula are an underestimate (Klug, 1958); while this is usually of little importance when considering the validity of a single sign determination it is more significant in this case where 86 probabilities are being multiplied together.

If the original scheme, when signs were developed by a step by step process with  $P<sub>L</sub>=0.80$ , is examined in the same way then it is found that some incorrect sign determinations are almost inevitable. With, say the 100th sign incorrectly determined experience suggests that the situation would soon deteriorate to the point where from about the 150th sign onwards completely random signs might be generated. Relationships between sign symbols deduced from the total resultant sign determination would be rather uncertain and there would thus be a high likelihood of failure.

This analysis also suggests an objective procedure as an alternative to initially fixing a value of *PL* by judgTable 2. Computer output for the development of signs for combinations of  $A, B, C = \pm 1$ 

The two most-probable combinations are shown.



ment. One may begin by finding new signs, one at a time, and calculating the cumulative product of probabilities at each step. When this has reached some value, 0.95 for example, all possible relationships between symbols can be found and if the probabilities of these are judged to be high enough then stage IV of the process can be commenced. Otherwise one goes further in finding new signs until the cumulative product has reached another limit, say 0.90. The search process is continued until the probability of the cumulative product is as high as possible compatible with having sufficiently high probabilities for the relationships between symbols. This procedure is not written into *LSAM* at the time of writing but may be included in due course.

#### **Non-centrosymmetric structures**

The multi-solution method of solving crystal structures proposed by Germain & Woolfson (1968) has been fully automated in the *MULTAN* computer program and a number of fairly complex structures (up to 46 atoms in the asymmetric unit,  $P2_12_12_1$ ) have been solved in this way.

Various trial procedures for phase development have been tried. They all had the general character that the strongest new phase indication from the tangent formula was accepted and incorporated in the known set and that after  $n$  such new phase determinations cyclic refinement of phases by the tangent formula was carried out. The various techniques tried were:

- (1) Taking  $n$  small, about five to ten, so that one always went forward with a highly self-consistent set of phases.
- (2) Taking n very large so that new phases were accepted in large numbers before refinement was carried out.
- (3) Making successive values of  $n$  increase in some exponential way, *i.e.* 5, 10, 20, 40 *etc.*

The trials showed, rather surprisingly, that frequent refinement was not only time-consuming but could be positively harmful. For example in one trial with space group  $P2_12_12_1$  and 142 reflexions the starting set contained seven reflexions. With  $n=10$  the final mean error in determined phase angle was 85.2°. For  $n=20$  or any higher value up to 135 the mean error in determined phase angle was about 30°!

However the fastest and most efficient process of all was found to be the use of a weighted tangent formula. Whenever a new phase,  $\varphi_h$ , is determined there is also calculated a corresponding weight,  $w_h$ . An effective scheme, incorporated in a routine called *FASTAN,*  was found to be

$$
\tan \varphi_{\mathbf{h}} \simeq \frac{\sum_{\mathbf{h}'} w_{\mathbf{h}'} w_{\mathbf{h} - \mathbf{h}'} | E_{\mathbf{h}'} E_{\mathbf{h} - \mathbf{h}'}| \sin (\varphi_{\mathbf{h}'} + \varphi_{\mathbf{h} - \mathbf{h}'})}{\sum_{\mathbf{h}'} w_{\mathbf{h}'} w_{\mathbf{h} - \mathbf{h}'} | E_{\mathbf{h}'} E_{\mathbf{h} - \mathbf{h}'}| \cos (\varphi_{\mathbf{h}'} + \varphi_{\mathbf{h} - \mathbf{h}'})} = \frac{T_{\mathbf{h}}}{B_{\mathbf{h}}} (3)
$$

$$
w_{\mathbf{h}} = \tanh \{ \sigma_3 \sigma_2^{-3/2} | E_{\mathbf{h}}| (T_{\mathbf{h}}^2 + B_{\mathbf{h}}^2)^{1/2} \} . \tag{4}
$$

where  $\sigma_2$  and  $\sigma_3$  have been defined by Karle & Karle (1966).

Since the tangent formula is weighted, all reflexions are included the whole of the time and there is no rejection criterion for a phase other than the weight being calculated as very small or zero. The form of weighting ensures that poorly determined phases have little effect in the determination of other phases, while the fact that all phases are included in the phase determination process leads to very efficient propagation of phase knowledge throughout the data set. It has been found that the weighted tangent formula converges on the final result significantly more quickly than the normal unweighted formula where one starts with a small subset of  $E$ 's and only after several cycles is the whole of the data included. In the authors' own case, the weighted tangent formula is about four times faster than the more conventional approach and on at least one occasion gave the correct structure when the unweighted tangent formula did not.

The weighting scheme, as proposed in equations (3) and (4), is an empirical one which is found to give sensible results. For this, as with most weighting schemes, the important thing is to get the right general form for the weighting function and the exact analytical expression used is normally not very critical.

The final phases as given by *FASTAN* have associated weights and a small weight indicates an uncertain phase. It can be shown by an approximate analysis that, taking these uncertainties into consideration, a Fourier synthesis with suitably weighted coefficients *KIEI* gives a higher signal: noise ratio than a normal  $|E|$  map.

Let us consider a Fourier map

$$
\Psi(\mathbf{r}) = \sum_{\mathbf{h}} K_{\mathbf{h}} |E_{\mathbf{h}}| \cos (2\pi \mathbf{h} \cdot \mathbf{r} - \varphi_{\mathbf{h}})
$$
 (5)

at  $r_j$ , the centre of the *j*th atom

$$
\Psi(\mathbf{r}_j) = \sum_{\mathbf{h}} K_{\mathbf{h}} |E_{\mathbf{h}}| \cos (2\pi \mathbf{h} \cdot \mathbf{r}_j - \varphi_{\mathbf{h}}) \,. \tag{6}
$$

It can be shown by arguments similar to those given by Cochran & Woolfson (1955) that the expectation value of cos  $(2\pi h \cdot r_i - \varphi_h)$  is

$$
\left\langle \cos\left(2\pi \mathbf{h} \cdot \mathbf{r}_j - \varphi_{\mathbf{h}}\right) \right\rangle = \frac{Z_j}{\left\{ \sum Z_i^2 \right\}^{1/2}} |E_{\mathbf{h}}| \tag{7}
$$

where  $Z_j$  is the atomic number of the *j*th atom and the summation is over all atoms in the unit cell. It is assumed in this that the form factors of the various atoms all have the same shape and differ only by a scale factor.

Replacing the cosine in expression (6) by its expectation value we have

$$
\langle \Psi(\mathbf{r}_j) \rangle = \frac{Z_j}{\{\sum Z_i^2\}^{1/2}} \sum_{\mathbf{h}} K_{\mathbf{h}} |E_{\mathbf{h}}|^2. \tag{8}
$$

If the errors in  $\varphi_h$  are not too large we can obtain an expression for the error in  $\Psi(\mathbf{r}_i)$  due to errors in the  $\varphi_{h}$ 's. This is

$$
\Delta \Psi(\mathbf{r}_j) = -\sum_{\mathbf{h}} K_{\mathbf{h}} |E_{\mathbf{h}}| \sin (2\pi \mathbf{h} \cdot \mathbf{r}_j - \varphi_{\mathbf{h}}) \Delta \varphi_{\mathbf{h}}. \tag{9}
$$

Since the  $\Delta\varphi_h$ 's are as equally likely to be positive as negative the expectation value of  $\Delta \Psi(\mathbf{r}_i)$  is zero.

However the mean-square value of  $\Delta \Psi(\mathbf{r}_1)$  is nonzero and equals

$$
\langle \langle \Delta \Psi(\mathbf{r}_j) \rangle^2 \rangle = \sum K_{\mathbf{h}}^2 |E_{\mathbf{h}}|^2 \langle \sin^2(2\pi \mathbf{h} \cdot \mathbf{r}_j - \varphi_{\mathbf{h}}) \Delta \varphi_{\mathbf{h}}^2 \rangle. \tag{10}
$$

Taking

$$
\langle \sin^2(2\pi \mathbf{h} \cdot \mathbf{r}_j - \varphi_{\mathbf{h}}) \rangle = 0.5
$$

and

$$
\langle \varDelta \varphi_{\mathbf{h}}^2 \rangle = V_{\mathbf{h}}
$$

the variance in the determination of  $\varphi_h$ , we find the rootmean-square value of  $\Delta \Psi(\mathbf{r}_i)$ 

$$
\langle \{ \Delta \Psi(\mathbf{r}_j) \}^2 \rangle^{1/2} = \{ \frac{1}{2} \sum_{\mathbf{h}} K_{\mathbf{h}}^2 | E_{\mathbf{h}} |^2 V_{\mathbf{h}} \}^{1/2} . \tag{11}
$$

For a high signal-to-noise ratio we might expect the ratio of the expression given in (8) to that given in (11) to be a maximum *i.e.* ignoring constant factors,

$$
X = \frac{\sum_{\mathbf{h}} K_{\mathbf{h}} |E_{\mathbf{h}}|^2}{\left\{\sum_{\mathbf{h}} K_{\mathbf{h}}^2 |E_{\mathbf{h}}|^2 V_{\mathbf{h}}\right\}^{1/2}}.
$$
 (12)

to be a maximum.

Multiplying all the K's by a constant factor does not affect the value of  $X$  so we can consider the problem as one of minimizing  $\Sigma K_h |E_h|^2 V_h$  while keeping  $\sum K_h |E_h|^2$  fixed.

Now

$$
d(\sum K_h |E_h|^2) = \sum_h |E_h|^2 dK_h = 0 \text{ for constant } \sum_h K_h |E_h|^2
$$

and

$$
d(\sum K_h^2 |E_h|^2 V_h) = 2 \sum_h K_h |E_h|^2 V_h dK_h
$$
  
= 0 for  $\sum_h K_h^2 |E_h|^2 V_h$  to be a minimum.

The values of  $K<sub>h</sub>$  satisfy the second condition if  $K_hV_h$  equals a constant.

It is found that if we interpret the quantity

$$
\alpha'_{\mathbf{h}} = 2\sigma_3 \sigma_2^{-3/2} |E_{\mathbf{h}}| (T_{\mathbf{h}}^2 + B_{\mathbf{h}}^2)^{1/2}
$$
 (13)

which is twice the argument of tanh in equation  $(4)$ , as corresponding to the normal  $\alpha_h$  as defined by Karle & Karle (1966) then  $K_h = \alpha_h$  is a reasonable approximation. This is shown in Table 3; the values of  $\alpha$  and V may be found tabulated in the paper by Germain, Main & Woolfson (1970).

Table 3. *Values of*  $\alpha$ *, V and*  $\alpha V$ 

V	$\alpha V$	α	V	α $V$
$(radians2) (radians2)$			$(radians2) (radians2)$	
3.290	0.000	6.0	0.184	1.104
2.349	1.175	8·0	0.134	1.072
1.604	1.604	$10-0$	0.106	1.060
0.764	1.528	$12 - 0$	0.087	1.044
0.437	1.311	$16 - 0$	0.065	1.040
0.298	1.192	20.0	0.051	1.020

The values of  $\alpha'_{h}$ , which vary between 0.6 and 56.9, are shown as ALPHA in Table 4 which is a typical output from *FASTAN* for a trial structure.

The above rough analysis suggests that  $\alpha'_{h}E_{h}$  should be used as a Fourier coefficient to get the greatest signal-to-noise ratio. However this analysis is rather rough and the derivation of equation (9) depends on small errors in the  $\varphi$ 's - which will not always be the case. The authors feel that weighting according to  $\alpha'_{h}$  is possibly too severe and that instead  $w_{h}$  may turn out to be a more sensible weighting function. For the extreme values of  $\alpha'_{h}$  of 0.6 and 56.9 the corresponding values of  $w_h$  are 0.291 and 1.000. Further investigations of weighting, both theoretical and empirical, will be made in due course.

#### **Additional information**

Germain, Main & Woolfson (1970) described the convergence method for selecting origin-and-enantiomorph-fixing reflexions together with some others to form a good starting set. As described in that paper the investigator chose origin-and-enantiomorph-fixing reflexions and these were tested for efficiency by the convergence method which also found the best other reflexions to take with them.

An improved version of the convergence method is incorporated in *MULTAN* as a routine called *CON-VERGE.* This does not require any intervention on the part of the investigator but automatically finds the best possible starting set of reflexions including those required to fix the origin and enantiomorph.

The routine *CONVERGE* also applies the  $\Sigma_1$  formula (Hauptman & Karle, 1953) to find probable signs of centric reflexions. The form of the  $\Sigma_1$  expressions depends on the space group; for  $P2_12_12_1$  for example one finds

$$
S\{E_{2h2k0}\}\simeq S\{\sum_{l}(-1)^{l+h}(|E^{hkl}|^2-1)\}\tag{14}
$$

with the probability that  $E_{2h,2k,0}$  is positive

$$
P_{+}\{E_{2h2k0}\} = \frac{1}{2} + \frac{1}{2} \tanh \left\{ \frac{1}{2} \sigma_{3} \sigma_{2}^{-3/2} | E_{2h,2k,0} | \right. \\ \times \sum_{l} (-1)^{l+h} (|E_{hkl}|^{2} - 1) \} . \tag{15}
$$

The weight associated with each sign determination is  $|2P-1|$  (see, for example, Woolfson, 1956) and  $\Sigma_1$ 

**phases with these weights can be incorporated into the starting sets. Their phases are kept constant and contribute to the phase determination with their initial weight until they are determined by the tangent formula with a larger weight. Thereafter they take on the value and weight indicated by the tangent formula. The other reflexions in the starting set are kept constant until the tangent formula refinement has essentially converged; they are then allowed to refine to their final values which takes about two further cycles.** 

**In Table 5 is shown the beginning and end of the output from** *CONVERGE.* **There were six signs found**  from  $\Sigma_1$  and these were all included in the starting set. **Although one of them (10,0,8) was incorrect the weighting scheme was able to handle the situation and a correct set of phases was found. There were two general phases whose phases were permuted in various**  combinations of  $\pm (\pi/4 \text{ or } 3\pi/4)$  giving sixteen different **starting points to be explored.** 

*MULTAN* **gives a number of possible sets of phases and in order to evaluate their relative plausibility it is necessary to have some figures of merit. One fairly obvious criterion is** 

$$
Z = \sum_{h} \alpha_{h} \tag{16}
$$

and it appears that it often enables the correct solution  
to be distinguished from incorrect solutions. However  
it is difficult to judge the worthiness of the values of  
Z in an absolute way – that is to say that one cannot  
tell whether or not the highest value of Z might  
conceivably correspond to a correct solution. For this  
reason we have developed an absolute figure of merit  
based on the statistics of the values of 
$$
\alpha
$$
. In German,  
Main & Woolfson (1970) it was shown that the ex-  
pected value of  $\alpha_h^2$ ,  $\langle \alpha_h^2 \rangle_e$ , is given by

$$
\langle \alpha_{\mathbf{h}}^2 \rangle_e = \sum_{\mathbf{h}'} K_{\mathbf{h},\mathbf{h}'}^2 + 2 \sum_{\substack{\mathbf{h}'} \to \mathbf{h}''} \sum_{\substack{\mathbf{h}''}} K_{\mathbf{h},\mathbf{h}'} K_{\mathbf{h},\mathbf{h}''} \frac{I_1(K_{\mathbf{h},\mathbf{h}'})}{I_0(K_{\mathbf{h},\mathbf{h}'})} \frac{I_1(K_{\mathbf{h},\mathbf{h}''})}{I_0(K_{\mathbf{h},\mathbf{h}''})}
$$

**where** 

$$
K_{\mathbf{h},\mathbf{h}'} = 2\sigma_3 \sigma_2^{-3/2} |E_{\mathbf{h}} E_{\mathbf{h}'} E_{\mathbf{h} - \mathbf{h}'}|.
$$

**For a random set of phases the expectation value of**   $\alpha_h^2$  is given by

$$
\langle \alpha_{\mathbf{h}}^2 \rangle_{\mathbf{r}} = \sum_{\mathbf{h}'} K_{\mathbf{h}, \mathbf{h}'}^2 \,. \tag{17}
$$

**One can thus use as an absolute figure of merit** 

$$
M_{\rm abs} = \frac{Z - \sum_{\rm h} \langle \alpha_{\rm h}^2 \rangle_r^{1/2}}{\sum \langle \alpha_{\rm h}^2 \rangle_c^{1/2} - \sum \langle \alpha_{\rm h}^2 \rangle_r^{1/2}} \,. \tag{18}
$$

# Table 4. *Values of*  $\alpha'_{h}$

**Computer output for** *FASTAN* **for a trial structure.** 

;"AFFINOSF FASTAN PHASF I}FIERHINATION II JUNE 1970

```
STARTING PHASES INPUT FROM CARDS
FOURIER OUTPUT NOT REQUIRED
```
EHp



# Table 5. Computer output for the computer program CONVERGE



 $\sim$ 

 $\sim 10^{11}$ 

The value of  $M_{\text{abs}}$  would tend to be near zero for random phases and of order unity for correct phases. In practice correct sets of phases are found to give values of  $M_{\text{abs}}$  in the range about 1.0-1.4. While some incorrect sets of phases can also give values of  $M<sub>abs</sub>$  of this magnitude it is unlikely that correct phases would give  $M_{\text{abs}}$  as low as 0.7, say. It should be stressed that no theoretical relationship has been found between  $M_{\text{abs}}$  and the probability of a set of phases being correct and that quoted judgments of the significance of values of  $M_{\text{abs}}$  are based on present  $experience - as yet limited.$ 

Table 4 shows the value of  $M$  as  $1.13$  for the correct set of phases for raffinose (Berman, 1970) a trial structure for trying out methods.

Provision is made in *FASTAN* for the input of starting phases from cards. This is particularly useful when a recycling procedure is being used in which a number of phases have been calculated from a partial structure and are to be used as a basis for further pbase determination. In this case it is advisable to fix most of the input phases for the initial stages of phase determination. This ensures that the new phases determined conform to the input information before refining all the phases together. If the input phases are allowed to refine immediately it is not uncommon for

them to change their values completely so they bear no relation to those input.

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# **Short-Range Ordering of Vacancies and Fermi Surface of TiO**

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The structure of the disordered (high-temperature) phase of titanium monoxide has been studied by electron diffraction. Composition-dependent complex distributions of diffuse scattering are observed and interpreted in terms of short-range-ordered configurations of lattice vacancies. The proposition that the form of the diffuse scattering may correspond to ordering under the influence of long-range electronic energy terms and so show a direct relationship to the topology of the Fermi surface is examined.

#### **Introduction**

The titanium monoxide phase has statistically the NaC1 type structure at high temperatures and is stable over a wide composition range *e.g.* from  $TiO<sub>0.70</sub>$  to  $TiO<sub>1.25</sub>$ at 1400 °C and from TiO<sub>0.90</sub> to TiO<sub>1.25</sub> at 900 °C and lower temperatures. The structure contains a large proportion of vacancies of both titanium and oxygen, the relative number of each type depending on composition. For  $TiO<sub>1.0</sub>$  15% of both kinds of site are vacant while  $TiO<sub>1-25</sub>$  has all oxygen sites full but 22% of the titanium sites empty.

Below a certain composition-dependent critical temperature, varying from 990 °C for  $TiO<sub>10</sub>$  to about 750 °C for TiO<sub>1.25</sub>, super-lattice structures are formed by the ordering of vacancies. These have been analysed by Watanabe, Castles, Jostsons & Malin (1967) and Watanabe, Terasaki, Jostsons & Castles (1968, 1969) on the basis of electron diffraction single-crystal patterns and by Hilti (1968) using X-ray diffraction patterns taken from crystals having a single orientation with respect to the basic NaCl-type structure but com-

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